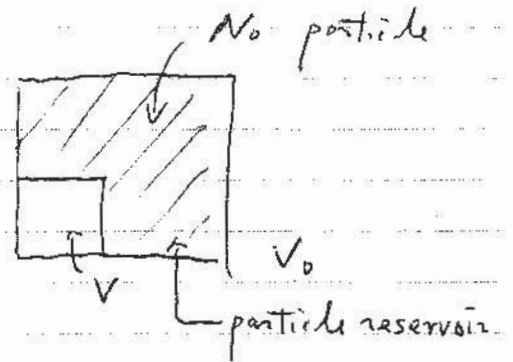


VS. Thermal fluctuations

① Thermal fluctuations of number of particle

Number of particle

$$\text{in } V: N = \sum_{i=1}^{N_0} x_i$$



$$x_i = 1 \quad \text{if } i\text{th particle is in } V$$

$$= 0 \quad \text{otherwise}$$

$$P(1) = \text{probability for } x=1 \quad P(1) = \frac{V}{V_0}$$

$$P(0) = \dots \quad \dots \quad x=0$$

$$\langle N \rangle = \sum_i \langle x_i \rangle = \sum_i P(1) = N_0 P(1)$$

$$\langle N^2 \rangle = \sum_i \langle x_i \rangle^2 + \sum_{i \neq j} \langle x_i x_j \rangle$$

$$= N_0 P(1) + (N_0^2 - N_0) \cdot P(1)^2 = \langle N \rangle + \langle N \rangle^2 - \langle N \rangle P(1)$$

$$\langle (N - \langle N \rangle)^2 \rangle \equiv \overline{\Delta N} = \langle N^2 \rangle - \langle N \rangle^2$$

$$= \langle N \rangle (1 - P(1)) \rightarrow \sqrt{\langle N \rangle} \quad \text{as } \frac{V}{V_0} \rightarrow 0$$

$$\boxed{\frac{\overline{\Delta N}}{\langle N \rangle} = \frac{1}{\sqrt{N}}}$$

* How to calculate the fluctuation of N for a generic system?

General ideas:

Thermal dynamical function S, A, Ω, G

⇒ Probability:

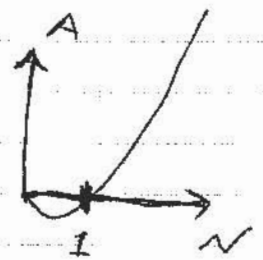
$$P \propto e^{S/k_B} \leftarrow \# \text{ of states} \quad P \propto e^{-\beta A} \leftarrow \text{total prob.}$$

$$P \propto e^{-\beta \Omega} \quad P \propto e^{-\beta G}$$

Gibbs potential

For ideal gas:

$$A(T, V, N) = N k_B T [\ln(N \lambda^3) - 1] \\ = N k_B T \ln N + \text{const}$$



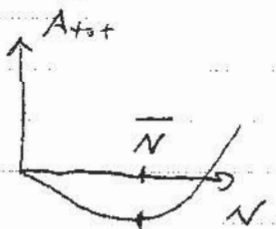
$P(N) = e^{-\beta A}$ is peaked at $N \approx \frac{1}{2}$

with particle reservoir: $P(N) = e^{-\beta A_{\text{tot}}}$

$$A_{\text{tot}}(T, V, N) = A - \mu N \\ = \text{const} + \frac{1}{2} \left. \frac{\partial^2 A}{\partial N^2} \right|_{\bar{N}} (N - \bar{N})^2$$

\bar{N} minimize A_{tot}

$$\frac{\partial A}{\partial N} = k_B T \ln N + k_B T, \quad \frac{\partial^2 A}{\partial N^2} = \frac{k_B T}{N}$$



$$P \propto e^{-\frac{\beta}{2} \frac{\partial^2 A}{\partial N^2} (N - \bar{N})^2} \Rightarrow \langle (N - \bar{N})^2 \rangle = \frac{k_B T}{\frac{\partial^2 A}{\partial N^2}} = \bar{N}$$

* What is $\frac{\partial^2 A}{\partial N^2} \Big|_{V,T} = \frac{\partial \mu}{\partial N} \Big|_{V,T}$? \Rightarrow compressibility

$$A(V, T, N) = N a(T, \frac{V}{N}) = \frac{V}{\sigma} a(T, \sigma)$$

$$\sigma = \frac{V}{N}$$

Change variable $(V, T, N) \Rightarrow (V, T, \sigma)$

$$\begin{aligned} \frac{\partial}{\partial N} \Big|_{V,T} &= \frac{\partial \sigma}{\partial N} \Big|_{V,T} \frac{\partial}{\partial \sigma} \Big|_{V,T} \\ &= -\frac{V}{N^2} \frac{\partial}{\partial \sigma} \Big|_{V,T} = -\frac{\sigma^2}{V} \frac{\partial}{\partial \sigma} \Big|_{V,T} \end{aligned}$$

$$\frac{\partial A}{\partial N} = -\frac{\sigma^2}{V} \frac{\partial}{\partial \sigma} \left(\frac{V}{\sigma} a(\sigma, T) \right) \Big|_{V,T}$$

$$\boxed{\frac{\partial a}{\partial \sigma} \Big|_T = \frac{\partial A}{\partial V} \Big|_{T,N} = P}$$

$$= a(\sigma, T) - \sigma \frac{\partial a}{\partial \sigma} \Big|_T = a(\sigma, T) - \sigma P(\sigma, T)$$

$$\frac{\partial^2 A}{\partial N^2} = -\frac{\sigma^2}{V} \frac{\partial}{\partial \sigma} [a(\sigma, T) - \sigma P(\sigma, T)] \Big|_{V,T}$$

$$= -\frac{\sigma^2}{V} \left[\frac{\partial a}{\partial \sigma} \Big|_T - P - \sigma \frac{\partial P}{\partial \sigma} \Big|_T \right] \quad \text{isothermal compressibility}$$

$$= \frac{\sigma^3}{V} \frac{\partial P}{\partial \sigma} = -\frac{\sigma^2}{V} \frac{1}{K_T}$$

$$\Rightarrow \boxed{\langle (N - \bar{N})^2 \rangle = k_B T \frac{V K_T}{\sigma^2}}$$

$$\boxed{K_T \equiv -\frac{1}{\sigma} \left(\frac{\partial V}{\partial P} \right) \Big|_T} \\ = -\frac{1}{V} \frac{\partial V}{\partial P} \Big|_T$$

$\langle (N - \bar{N})^2 \rangle \Rightarrow$ fluctuation

$K_T = -\frac{1}{V} \frac{\partial V}{\partial P} \Rightarrow$ response

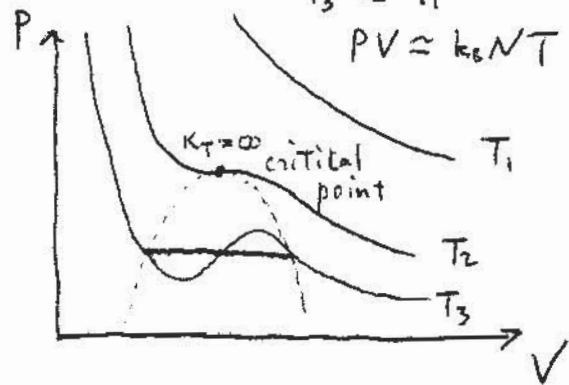
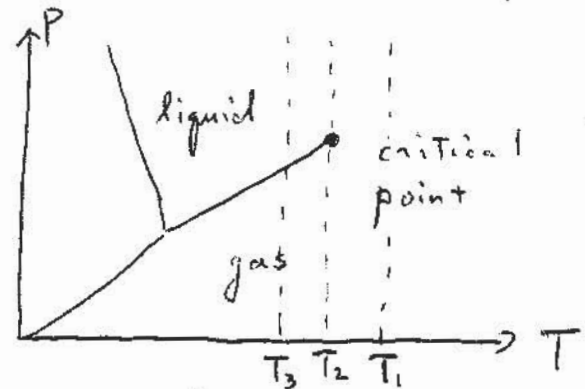
$\} \text{ closely related.}$

$$\Leftrightarrow \frac{\partial^2 A}{\partial N^2}$$

* δN has large fluctuations near the critical point

$$\begin{aligned}\frac{\delta n}{n} &= \frac{\sqrt{\langle (N - \bar{N})^2 \rangle}}{\bar{N}} \\ &= \sqrt{k_B T \frac{K_T}{V}} \frac{1}{\bar{N}} \\ &= \sqrt{k_B T K_T} \frac{1}{\bar{V}} \\ &= \sqrt{\frac{V_c}{V}}\end{aligned}$$

$$V_c = k_B T K_T$$



\Rightarrow Water density in a smaller volume has a stronger fluctuation
The fluctuations become of order 1 when $V \sim V_c$

Usually $V_c \sim (1 \text{ \AA})^3$

$V_c \rightarrow \infty$ at the critical point

When $V_c \sim (500 \text{ \AA})^3$, the water turn milky

(2) Fluctuations of V :

$$\star e^{-\beta(A(V,T) + VP)} = \# e^{-\frac{1}{2}\beta \frac{\partial^2 A}{\partial V^2} (V - \bar{V})^2}$$

$$\frac{\partial^2 A}{\partial V^2} \Big|_T = - \frac{\partial P}{\partial V} \Big|_T = - \frac{1}{\frac{\partial V}{\partial P} \Big|_T} = \frac{1}{VK_T}$$

K_T : compressibility
 $= \frac{1}{V} \frac{\partial V}{\partial P} \Big|_T$

For ideal gas:

$$V = \frac{Nk_B T}{P}, \quad VK_T = + \frac{Nk_B T}{P^2} = \frac{V}{P} = \frac{V^2}{Nk_B T}$$

$\langle (V - \bar{V})^2 \rangle$
 $= \frac{k_B T}{\frac{\partial^2 A}{\partial V^2}} = k_B T K_T V = \frac{V^2}{N}$

ideal gas

$\delta V = \sqrt{\langle (V - \bar{V})^2 \rangle} = \frac{V}{\sqrt{N}}$

ideal gas

\star Second way:

$$\frac{\partial}{\partial P} G = -k_B T \frac{\int dV (-\beta)V e^{-\beta[A(V,T) + VP]}}{\int dV e^{-\beta[A(V,T) + VP]}}$$

$$= + \langle V \rangle$$

$$\frac{\partial^2 G}{\partial P^2} = -\beta \frac{\int dV V^2 e^{-\beta(A+VP)}}{\int dV e^{-\beta(A+VP)}} - \frac{-\beta \left(\int dV V e^{-\beta(A+VP)} \right)^2}{\left(\int dV e^{-\beta(A+VP)} \right)^2}$$

$$= [\langle V^2 \rangle - \langle V \rangle^2] (-\beta)$$

$\langle (V - \bar{V})^2 \rangle = -k_B T \frac{\partial^2 G}{\partial P^2} = k_B T \cdot VK_T$

$\frac{\partial G(P,T)}{\partial P} = V$
 $P = - \frac{\partial A}{\partial V}$
 $V = \frac{\partial G}{\partial P}$
 $V \leftrightarrow P$ conjugate
 $A(V,T)$
 $G(P,V)$

Total entropy $S_{\text{tot}} = S(E) - \frac{E}{T}$ entropy of heat bath 46

sys

←

heat bath

③ Fluctuation of energies

$$\text{Prob.} \propto e^{-\frac{1}{2} \left[\frac{d^2}{dE^2} \frac{1}{k_B} \left[S(E) - \frac{E}{T} \right] \right]} (E - \bar{E})^2$$

$$= e^{-\frac{1}{2} k_B^{-1} \frac{d^2 S(E)}{dE^2} (E - \bar{E})^2}$$

$$\Rightarrow \langle (E - \bar{E})^2 \rangle = - \frac{k_B}{d^2 S(E) / dE^2}$$

$$\frac{\partial^2 S}{\partial E^2} = \frac{1}{T^2} \frac{\partial T}{\partial E}$$

$$= \frac{1}{T^2 C_V}$$

$$= -N \frac{k_B}{d^2 S / dE^2} = k_B T^2 C_V$$

$S = S/N \sim S/V$

$\epsilon = E/N \sim E/V$

$$\Rightarrow \delta E = \sqrt{\langle (E - \bar{E})^2 \rangle} \approx \sqrt{k_B T^2 C_V} \sim \sqrt{N}$$

$$\bar{E} \sim N \Rightarrow \bar{E} \gg \sqrt{\langle (E - \bar{E})^2 \rangle}$$

small fluctuation

$$\frac{\delta E}{\bar{E}} \sim \frac{1}{\sqrt{N}}$$

* For ideal gas $S = k_B \ln \frac{V (2mE)^{3/2}}{h^3}$

$$\frac{d^2 S}{dE^2} = k_B \frac{3}{2} \frac{1}{E^2}$$

$$\delta E = \sqrt{N} \sqrt{\frac{2}{3}} \epsilon = \sqrt{\frac{2}{3}} \frac{E}{\sqrt{N}}$$

④ Fluctuation and response

In general

$$A(x)$$

↑ some parameter

which can fluctuate

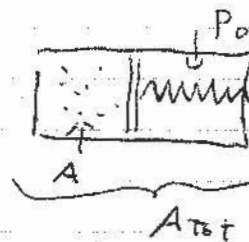
$$P \propto e^{-\beta A(x)}$$

$$A = \text{const} + \frac{1}{2} \frac{\partial^2 A}{\partial x^2} (x - \bar{x})^2$$

$$\Rightarrow \langle (x - \bar{x})^2 \rangle = \frac{k_B T}{\frac{\partial^2 A}{\partial x^2} |_{\bar{x}}}$$

$$A_{\text{tot}} = A + P_0 V$$

$$\langle (V - \bar{V})^2 \rangle = \frac{k_B T}{\frac{\partial^2 A_{\text{tot}}}{\partial V^2}}$$



$$= k_B T \frac{\partial^2 A}{\partial V^2} = k_B T \left(-\frac{\partial P}{\partial V} \right)_T$$

Ideal gas

$$PV = k_B T N$$

$$-\frac{\partial P}{\partial V} = \frac{k_B T N}{V^2}$$

$$\Rightarrow \langle (V - \bar{V})^2 \rangle = \frac{\bar{V}^2}{N} = \Delta V^2$$

$$\Delta V = \frac{V}{N}$$

$$\frac{\Delta V}{V} = \frac{1}{N}$$

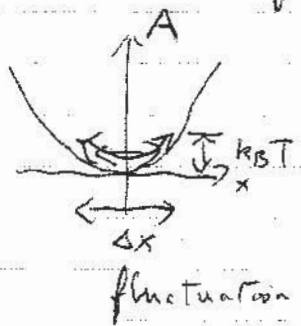
$$\langle (V - \bar{V})^2 \rangle = \frac{k_B T}{\left(-\frac{\partial P}{\partial V} \right)_T} = k_B T \left(-\frac{\partial V}{\partial P} \right)_T$$

↑ fluctuation

↑ response

$$\frac{\partial V}{\partial P} \quad P \leftrightarrow \text{"force"}$$

Thermal fluctuations - response



$$\langle \Delta x^2 \rangle = k_B T K$$

