

PROFESSOR: Today we'll talk about observables and Hermitian operators. So we've said that an operator, Q , is Hermitian in the language that we've been working so far, if you find that the integral, $\int \psi_1^* Q \psi_2 dx$, is actually equal to the integral $\int dx$ of Q , acting this time of ψ_1 all star ψ_2 .

So as you've learned already, this requires some properties about the way functions far away, at infinity, some integration by parts, some things to manage, but this is the general statement for a large class of functions, this should be true. Now we want to, sometimes, use a briefer notation for all of this. And I will sometimes use it, sometimes not, and you do whatever you feel. If you like to use this notation, use it. So here's the definition. If you put up ψ_1 , ψ_2 and a parentheses, this denotes a number, and in fact denotes the integral of ψ_1^* of x , ψ_2 of $x dx$.

So whatever you put in the first input ends up complex conjugated. When you put in the second input, it's like that, it's all integrated. This has a couple of obvious properties. If you put a number times ψ_1 times ψ_2 like this, the number will appear, together with ψ_1 , and will complex conjugated. So it can go out as a star $\psi_1 \psi_2$. And if you put the number on the second input, it comes out as is. Because the second input is not complex conjugated in the definition. With this definition, a Hermitian operator, Q is Hermitian, has a nice look to it. It becomes kind of natural and simple.

It's the statement that if you have ψ_1 , $Q \psi_2$, you can put the Q in the first input. $\int \psi_1^* Q \psi_2 dx$. This second term in the right hand side is exactly this integral here. And the first term in the left hand side is the left hand side of that condition. So it's just maybe a briefer way to write it. So when you get tired of writing $\int dx$ of the first, the second, you can use this.

Now with distance last time, the expectation values of operators. So what's the expectation value of Q in some state ψ of x ? And that is denoted as these braces here and of ψ is equal to the integral of ψ^* . The expectation value depends on the state you live in and it's $\int \psi^* Q \psi dx$. Or if you wish, in written notation $\langle \psi | Q | \psi \rangle$. I should put the hats everywhere. This is the expectation value of Q . I'm sorry, I missed here a star. So so far, so good. We've reviewed what a Hermitian operator is, what an expectation value is, so let's begin with some claims.

Claim number one. The expectation value of Q , with Q Hermitian. So everywhere here, Q will be Hermitian. The expectation value of Q is real. A real number, it belongs to the real

numbers. So that's an important thing. You want to figure out the expectation value of Q , you have a ψ^* , you have a ψ . Well, it'd better be real if we're going to think, and that's the goal of this discussion, that Hermitian operators are the things you can measure in quantum mechanics, so this better be real.

So let's see what this is. Well, $\langle Q \rangle_{\psi}$, that's the expectation value. If I complex conjugate it, I must complex conjugate this whole thing. Now if you want to complex conjugate an integral, you can complex conjugate the integrand. Here it is. I took this right hand side here, the integrand. I copied it, and now I complex conjugated it. That's what you mean by complex conjugating an integral. But this is equal, $\int dx$. Now I have a product of two functions here. ψ^* and Q that has acted on ψ . So that's how I think. I never think of conjugating Q . Q is a set of operations that have acted on ψ and I'm just going to conjugate it. And the nice thing is that you never have to think of what is Q^* , there's no meaning for it.

So what happens here? Priority of two functions, the complex conjugate of the first-- now if you [INAUDIBLE] normally something twice, you get the function back. And here you've got $\langle Q \rangle_{\psi}$. But that, these are functions. You can move around. So this $\langle Q \rangle_{\psi}^* = \langle \psi | Q | \psi \rangle$. And so far so good. You know, I've done everything I could have done. They told to come to complex conjugate this, so I complex conjugated it and I'm still not there. But I haven't used that this operator is Hermitian. So because the operator is Hermitian, now you can move the Q from this first input to the second one. So it's equal to $\int dx \psi^* Q \psi$. And oh, that was the expectation value of Q on ψ , so the star of this number is equal to the number itself, and that proves the claim, Q is real.

So this is our first claim. The second claim that is equally important, claim two. The eigenvalues of the operator Q are real. So what are the eigenvalues of Q ? Well you've learned, with the momentum operator, eigenvalues or eigenfunctions of an operator are those special functions that the operator acts on them and gives you a number called the eigenvalue times that function. So $Q \psi_1 = \lambda \psi_1$, if ψ_1 is a particularly nice choice, then it will be equal to some number. Let me quote $\lambda \psi_1$. And there, I will say that λ is the eigenvalue. That's the definition. And ψ_1 is the eigenvector, or the eigenfunction. And the claim is that that number is going to be real.

So why would that be the case? Well, we can prove it in many ways, but we can prove it kind of easily with claim number one. And actually gain a little insight, could calculate the expectation value of Q on that precise state, ψ_1 . Let's see how much is it. You see, ψ_1 is a particular

state. We've called it an eigenstate of the operator. Now you can ask, suppose you live in ψ_1 ? That's who you are, that's your state. What is the expectation value of this operator? So we'll learn more about this question later, but we can just do it, it's the integral of $\psi_1^* Q \psi_1$. And I keep forgetting these stars, but I remember them after a little while. So at this moment, we can use the eigenvalue condition, this condition here, that this is equal to $\psi_1^* Q \psi_1$. And the $Q \psi_1$ can go out, hence $Q \int \psi_1^* \psi_1$.

But now, we've proven, in claim number one, that the expectation value of Q is always real, whatever state you take. So it must be real if you take it on the state ψ_1 . And if the expectation value of ψ_1 is real, then this quantity, which is equal to that expectation value, must be real. This quantity is the product of two factors. A real factor here-- that integral is not only real, it's even positive-- times $Q \psi_1$. So if this is real, then because this part is real, the other number must be real. Therefore, $Q \psi_1$ is real.

Now it's an interesting observation that if your eigenstate, eigenfunction is a normalized eigenfunction, look at the eigenfunction equation. It doesn't depend on what precise ψ_1 you have, because if you put ψ_1 or you put twice ψ_1 , this equation still holds. So if it holds for ψ_1 , if ψ_1 is called an ideal function, $3\psi_1$, $5\psi_1$, ψ_1 are all eigenfunctions.

Properly speaking in mathematics, one says that the eigenfunction is the subspace generated by this thing, by multiplication. Because everything is accepted. But when we talk about the particle maybe being in the state of ψ_1 , we would want to normalize it, to make $\int \psi_1^* \psi_1$ squared equal to 1. In that case, you would obtain that the expectation value of the operator on that state is precisely the eigenvalue. When you keep measuring this operator, this state, you keep getting the eigenvalue. So I'll think about the common for a normalized ψ_1 as a true state that you use for expectation values.

In fact, whenever we compute expectation values, here is probably a very important thing. Whenever you compute an expectation value, you'd better normalize the state, because otherwise, think of the expectation value. If you don't normalize the state, you the calculation and you get some answer, but your friend uses a wave function three times yours and your friend gets now nine times your answer. So for this to be a well-defined calculation, the state must be normalized.

So here, we should really say that the state is normalized. Say one is the ideal function normalized. And this integral would be equal to $Q \psi_1$ belonging to the reals. And $Q \psi_1$ is real. So

for a normalized ψ_1 or how it should be, the expectation value of Q on that eigenstate is precisely equal to the eigenvalue.