

THE SHELL MODEL

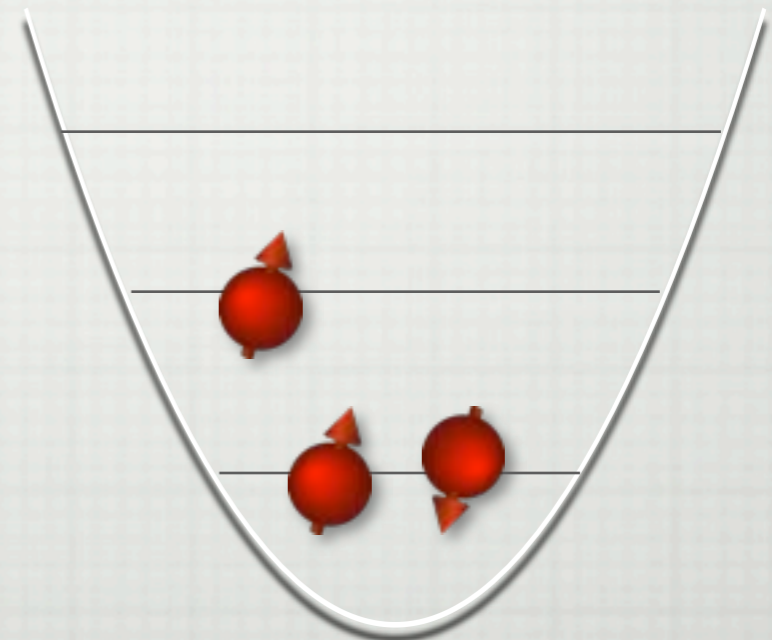
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Introduction To Applied Nuclear Physics

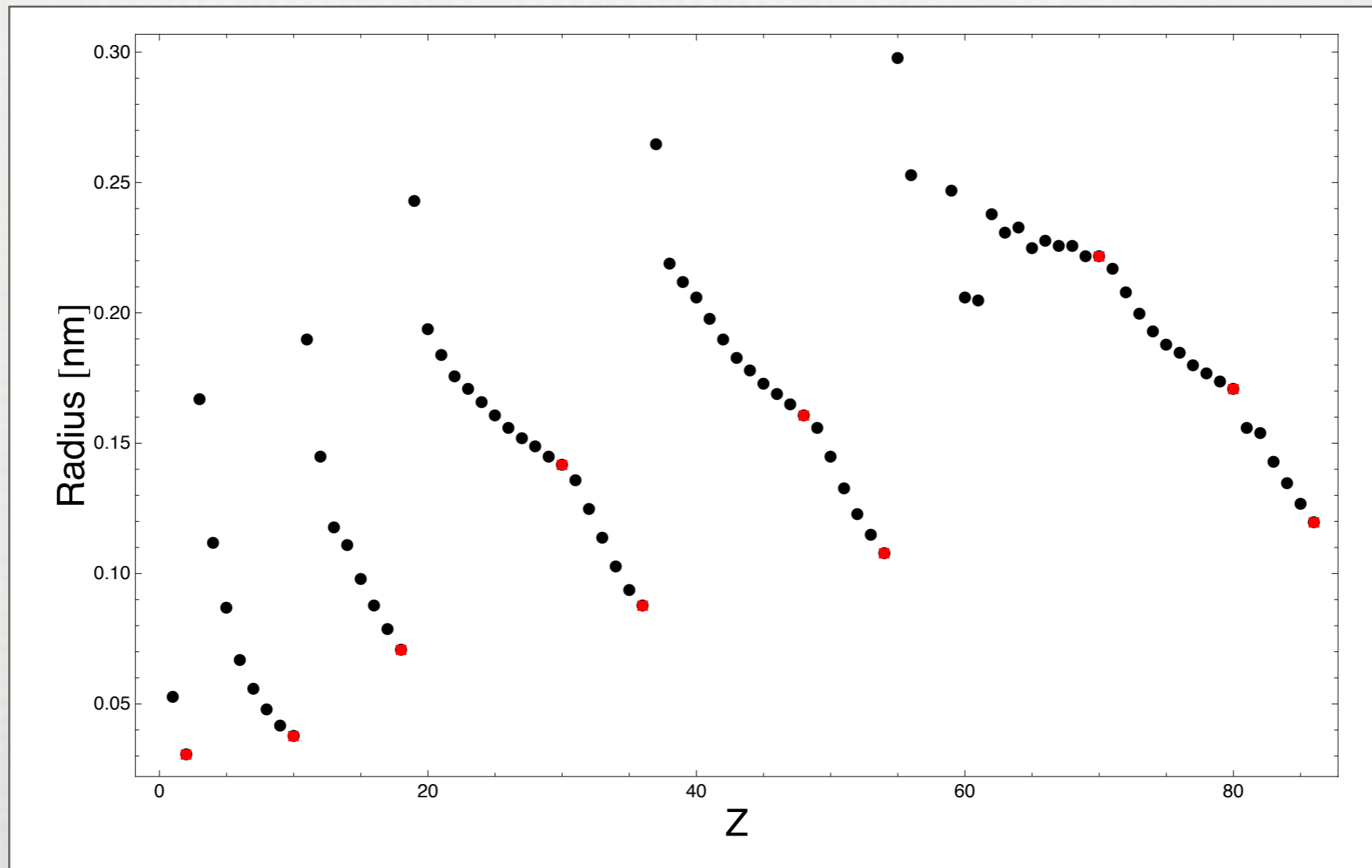
Spring 2012

Atomic Shell Model

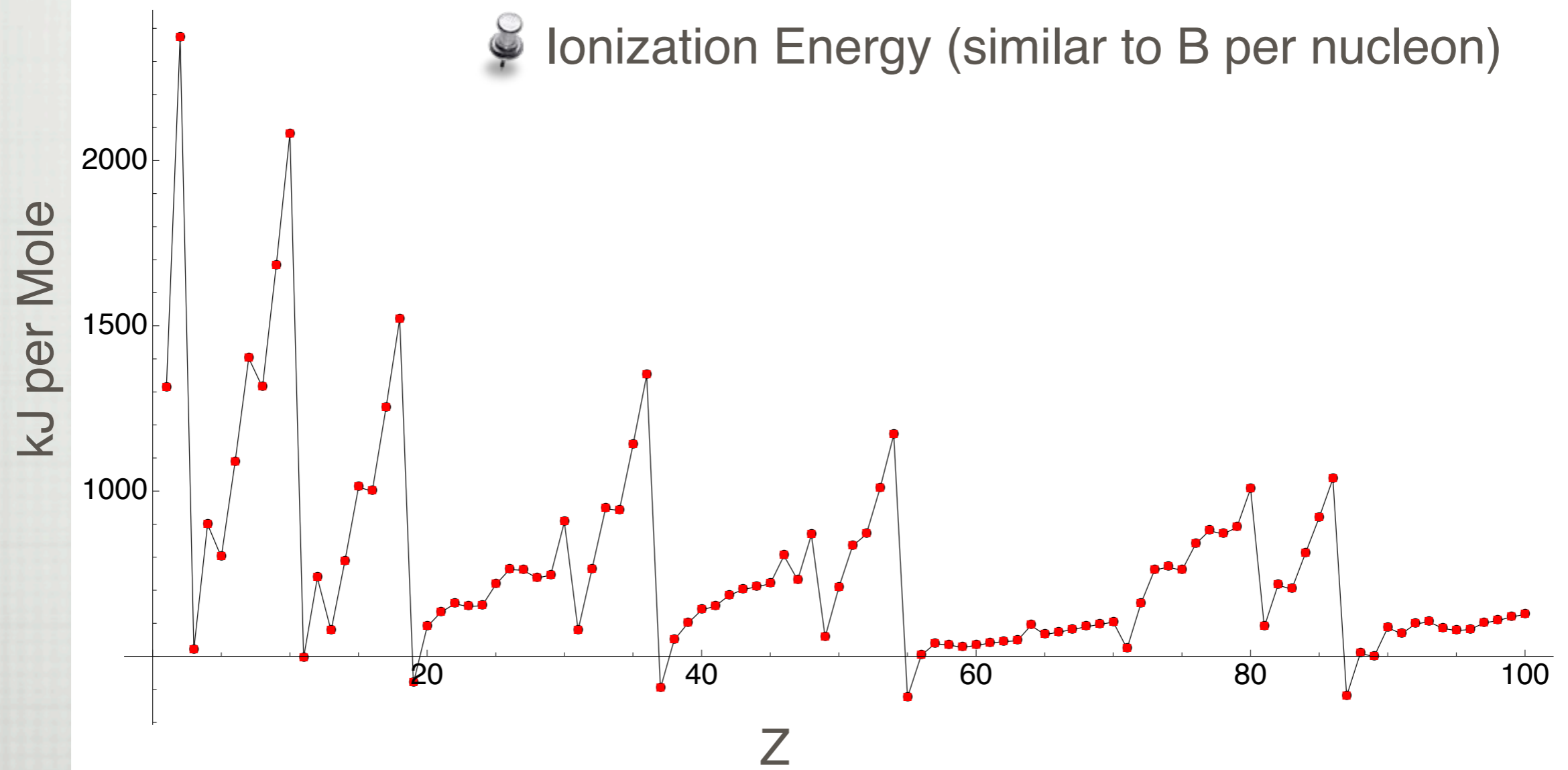
- Chemical properties show a periodicity
- Periodic table of the elements
- Add electrons into shell structure



Atomic Radius



Ionization Energy



ATOMIC STRUCTURE

• The atomic wavefunction is written as

$$|\psi\rangle = |n, l, m\rangle = R_{n,l}(r)Y_l^m(\vartheta, \varphi)$$

• where the labels indicate :

- n : principal quantum number
- l : orbital (or azimuthal) quantum number
- m : magnetic quantum number

• The degeneracy is

$$\mathcal{D}(l) = 2(2l + 1) \rightarrow \mathcal{D}(n) = 2n^2$$

AUFBAU PRINCIPLE

 The orbitals (or shells) are then given by the n-levels (?)


l	0	1	2	3	4	5	6
Spectroscopic notation	s	p	d	f	g	h	i
$\mathcal{D}(l)$	2	6	10	14	18	22	26
	historic structure				heavy nuclei		

n	$\mathcal{D}(n)$	e^- in shell
1	2	2
2	6	8
3	18	28

ATOMIC PERIODIC TABLE

H																	He
Li	Be											B	C	N	O	F	Ne
Na	Mg											Al	Si	P	S	Cl	Ar
K	Ca	Sc	Ti	V	Cr	Mn	Fe	Co	Ni	Cu	Zn	Ga	Ge	As	Se	Br	Kr
Rb	Sr	Y	Zr	Nb	Mo	Tc	Ru	Rh	Pd	Ag	Cd	In	Sn	Sb	Te	I	Xe
Cs	Ba	Lu	Hf	Ta	W	Re	Os	Ir	Pt	Au	Hg	Tl	Pb	Bi	Po	At	Rn
Fr	Ra																

AUFBAU PRINCIPLE

 The orbitals (or shells) are then given by close-by energy-levels

l	0	1	2	3	4	5	6
Spectroscopic notation	s	p	d	f	g	h	i
$\mathcal{D}(l)$	2	6	10	14	18	22	26
	historic structure				heavy nuclei		

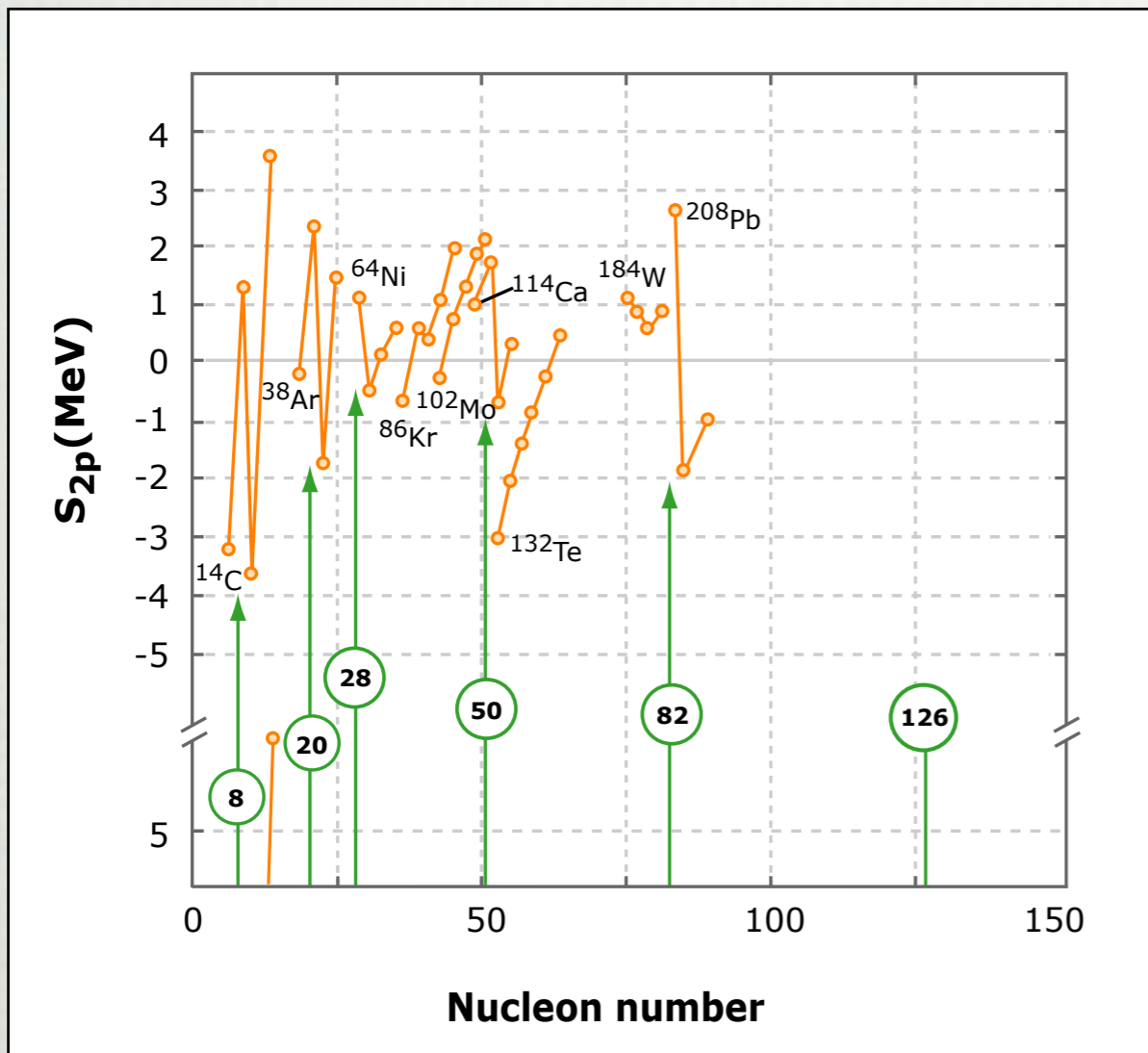
n	$\mathcal{D}(n)$	e^- in shell
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3s+3p form one level with # 10
4s is filled before 3d

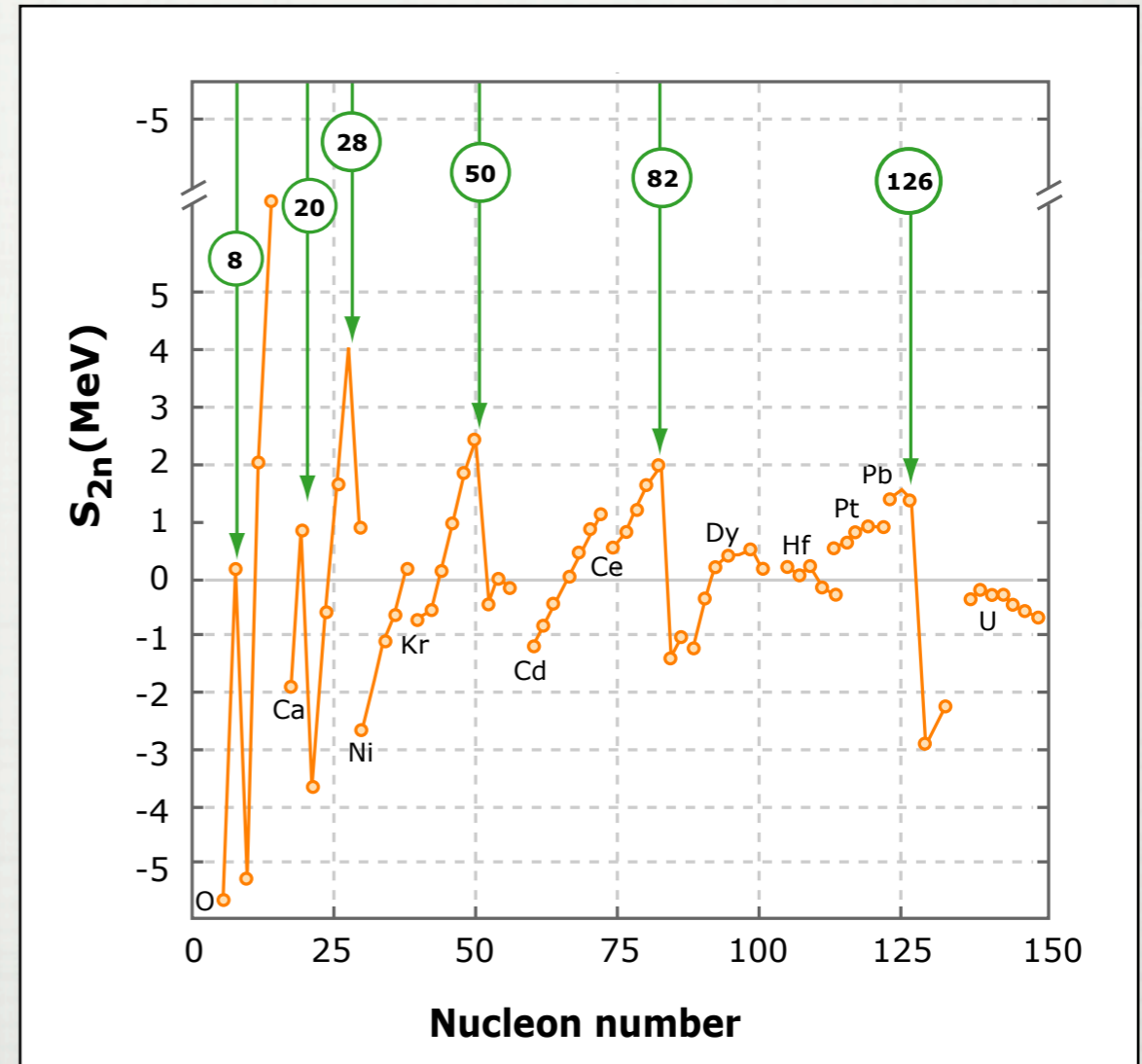
Nuclear Shell Model

- Picture of adding particles to an external potential is no longer good: each nucleon contributes to the potential
- Still many evidences of a shell structure

Separation Energy

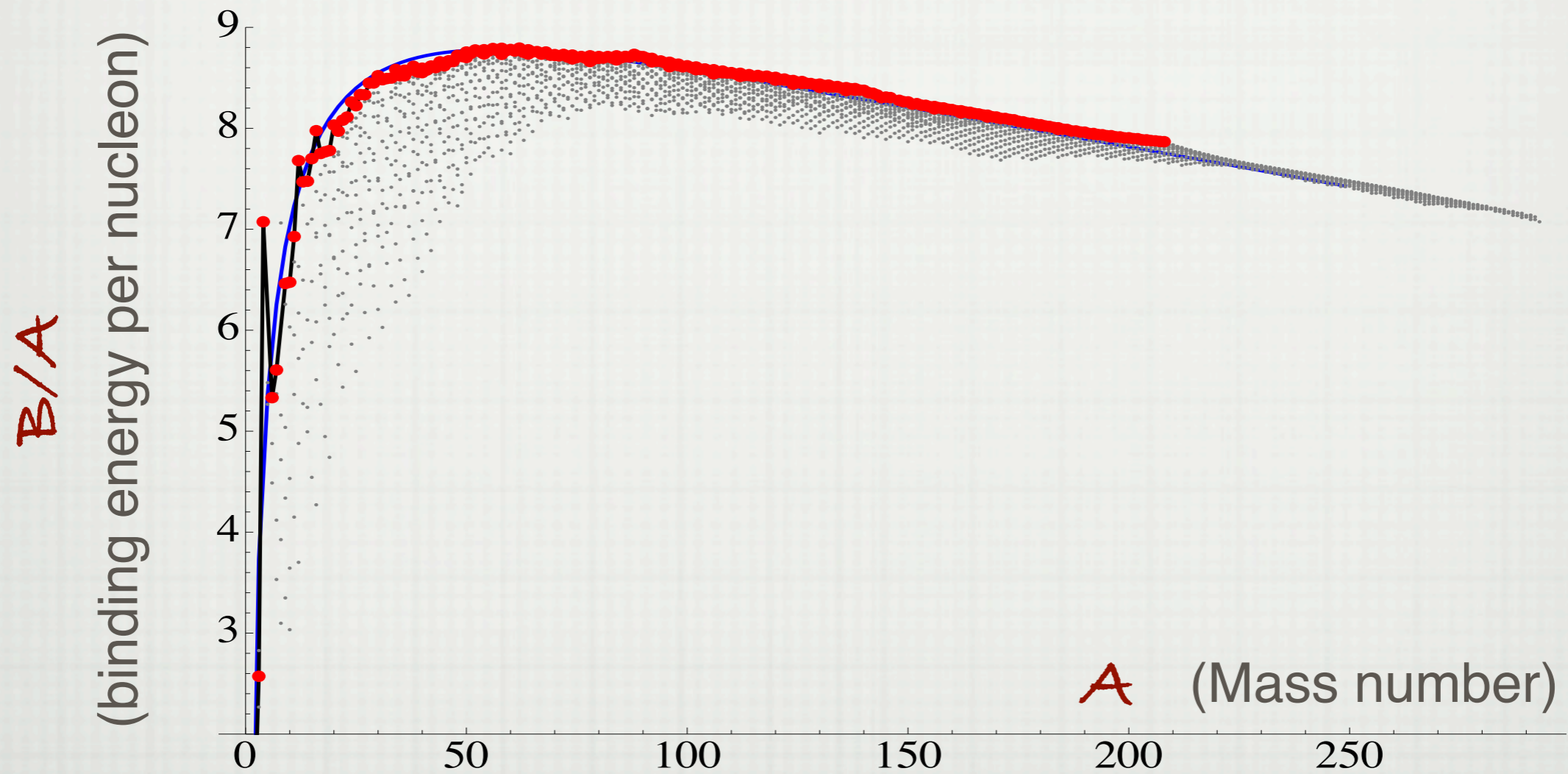


PROTON



NEUTRON

B/A: JUMPS



● "Jumps" in Binding energy from experimental data

CHART of NUCLIDES (Z/A vs. A)

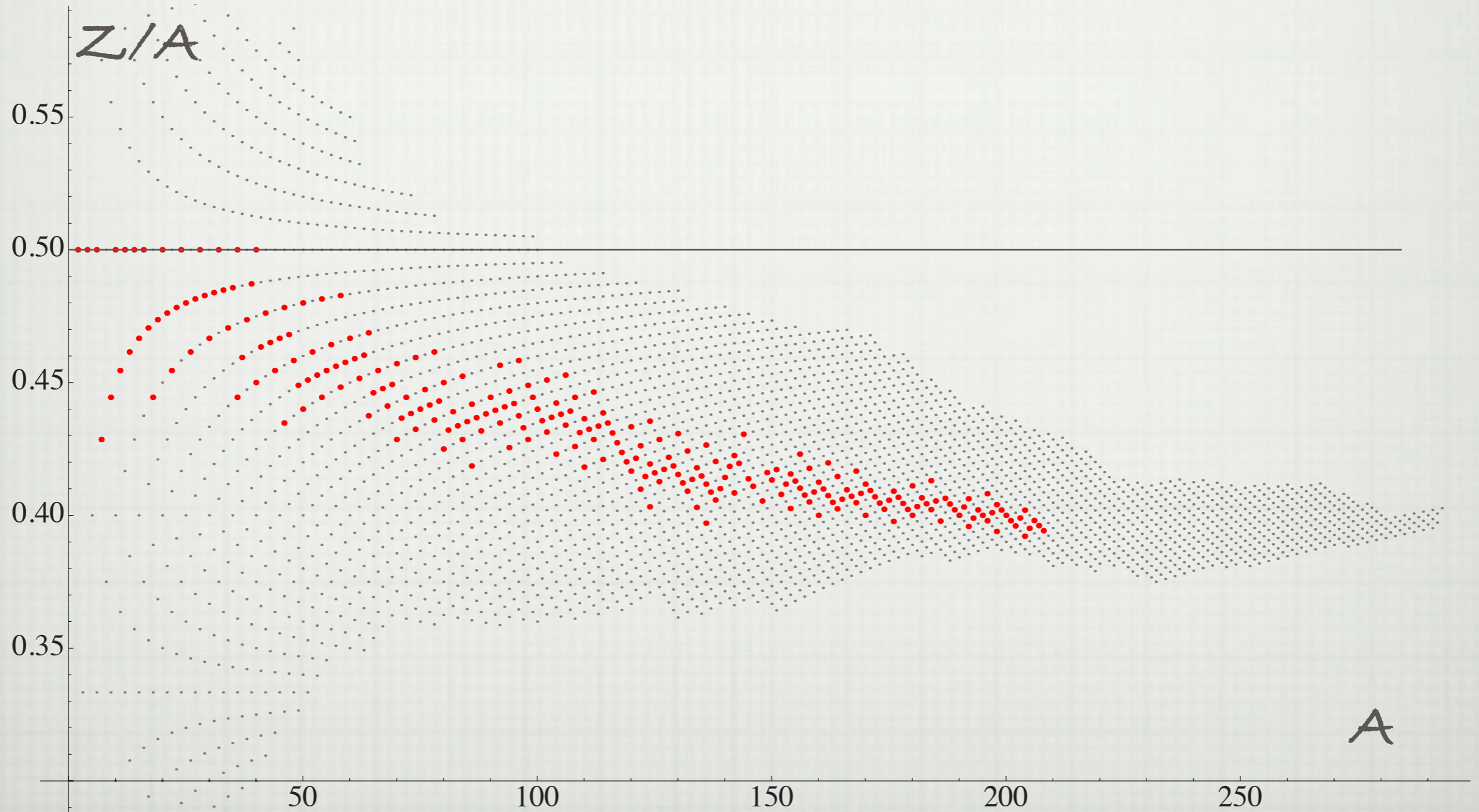


CHART OF NUCLIDES



“Periodic”, more complex properties → nuclear structure

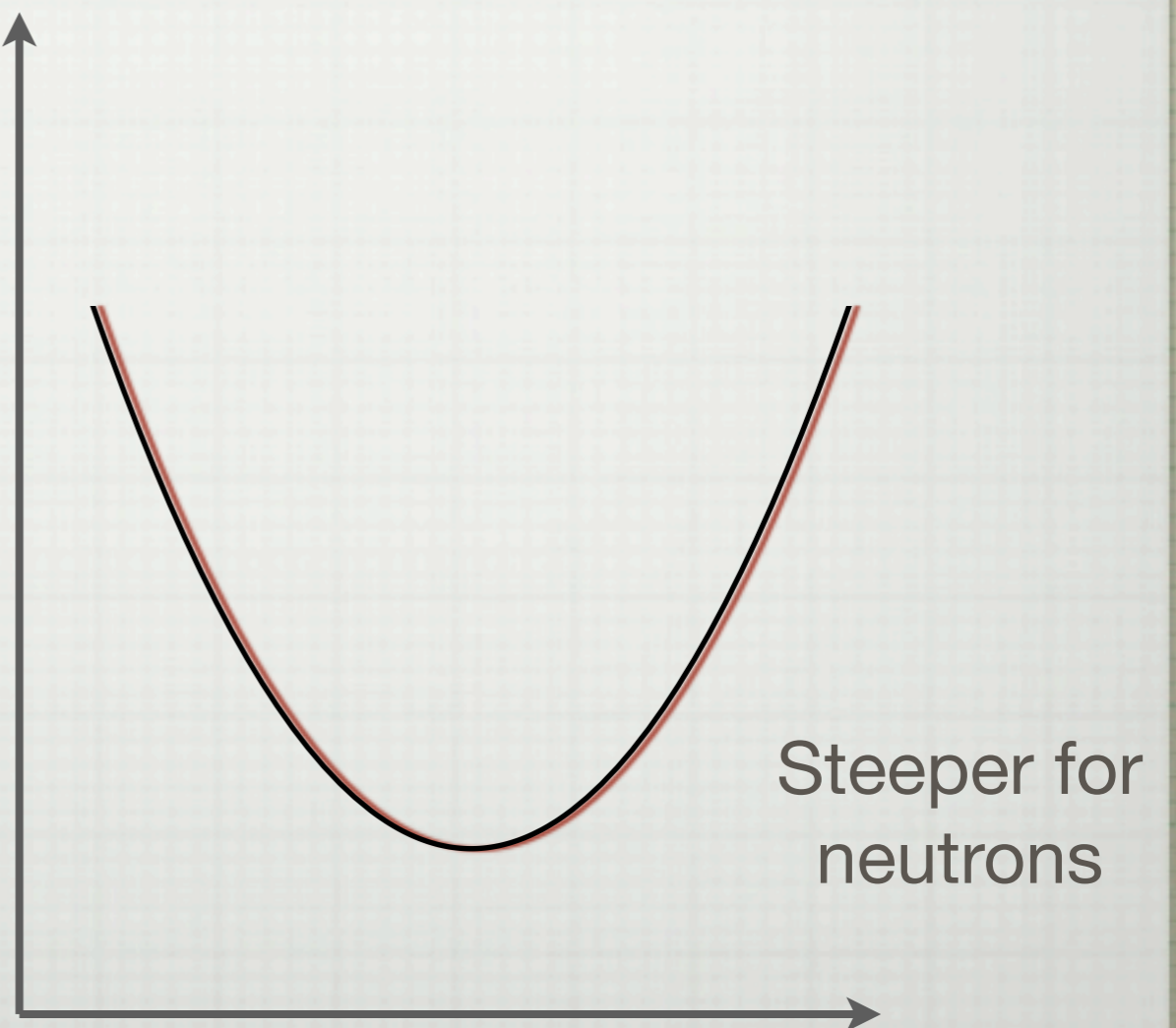
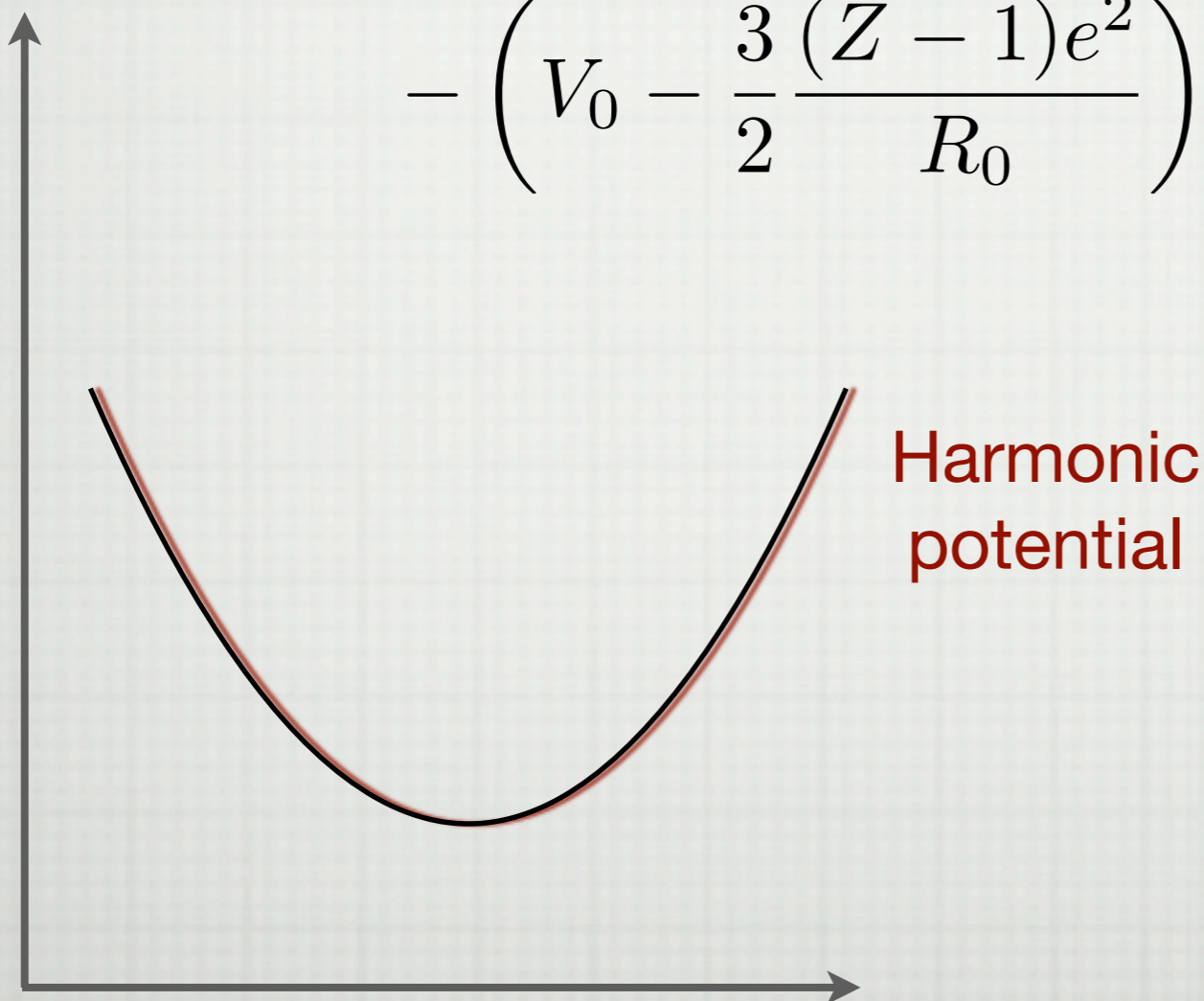
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NUCLEAR POTENTIAL

$$V_p = r^2 \left(\frac{V_0}{R_0^2} - \frac{(Z-1)e^2}{2R_0^3} \right)$$

$$V_n = r^2 \left(\frac{V_0}{R_0^2} \right) - (V_0)$$

$$- \left(V_0 - \frac{3}{2} \frac{(Z-1)e^2}{R_0} \right)$$



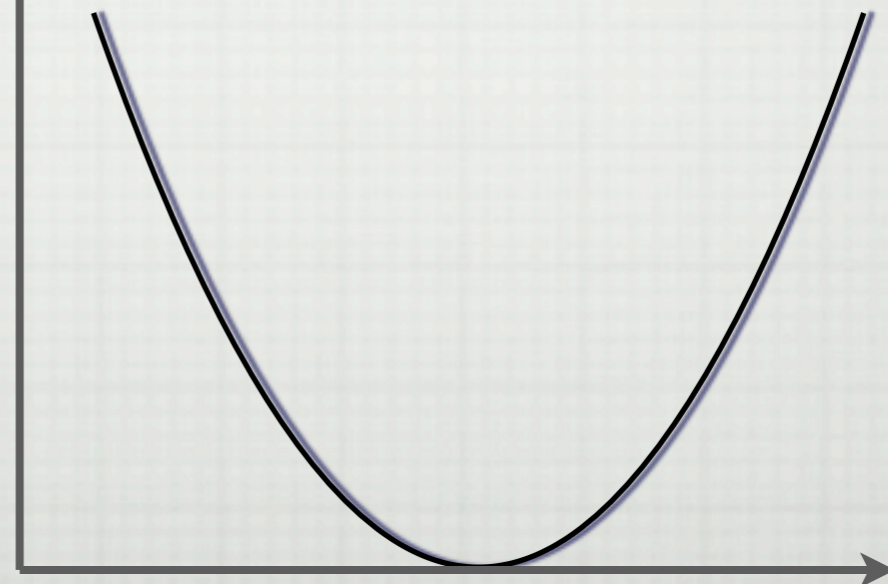
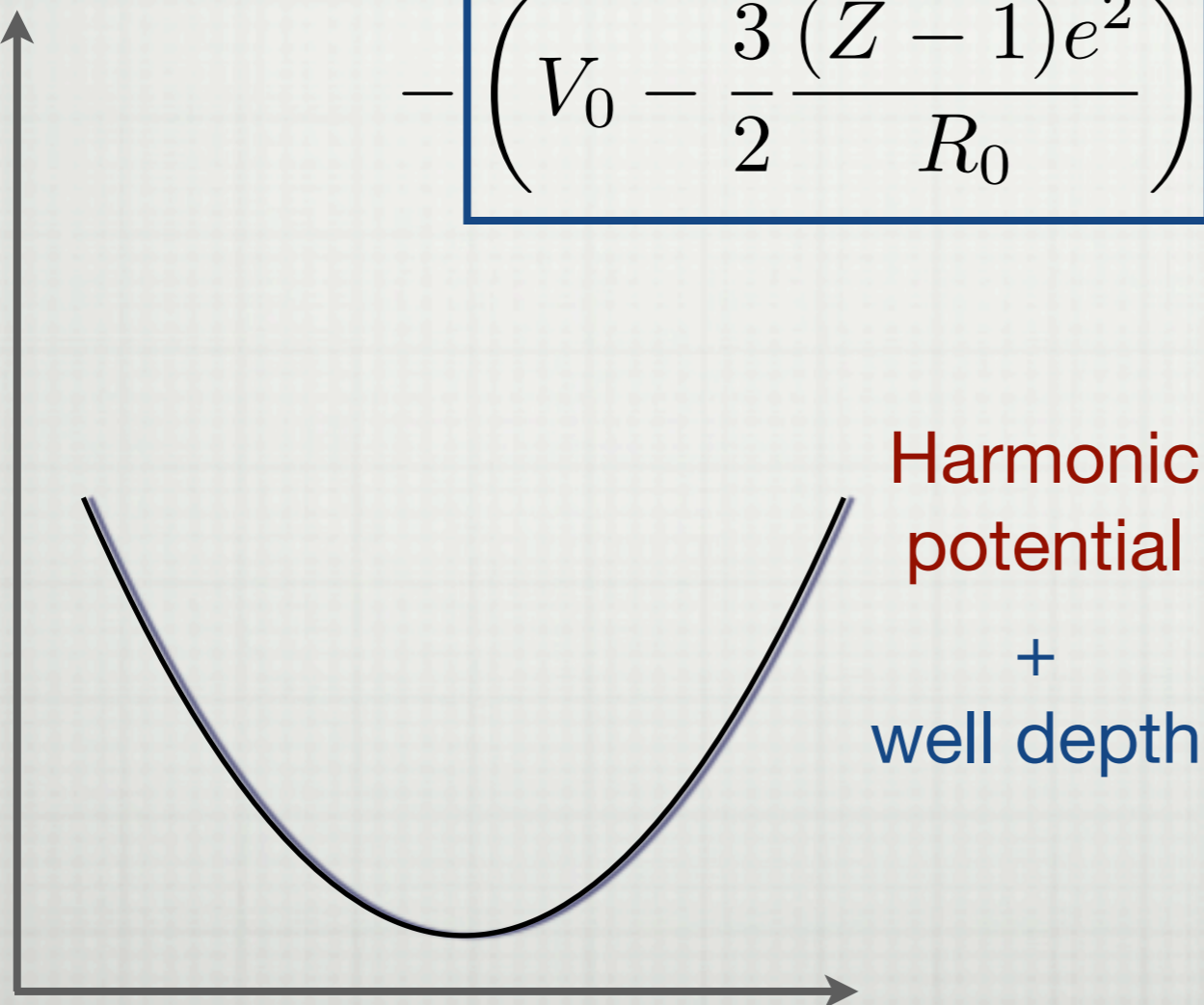
NUCLEAR POTENTIAL

$$V_p = r^2 \left(\frac{V_0}{R_0^2} - \frac{(Z-1)e^2}{2R_0^3} \right)$$

$$V_n = r^2 \left(\frac{V_0}{R_0^2} \right) - (V_0)$$

$$- \left(V_0 - \frac{3}{2} \frac{(Z-1)e^2}{R_0} \right)$$

Steeper and Deeper
for neutrons



Shell Mode

- 📌 Harmonic oscillator: solve (part of) the radial equation
 - including the angular momentum (centrifugal force term) we obtain the usual principal quantum number $n = (N-1)/2+1$

N	l	Spectroscopic Notation	$\mathcal{D}(N)$	$2\mathcal{D}(N)$	Cumulative #
0	0	1 s	1	2	2
1	1	1 p	3	6	8
2	0,2	2s,1d	6	12	20
3	1,3	2p,1f	10	20	40
4	0,2,4	3s,2d,1g	15	30	70

Spin-Orbit Coupling

- The spin-orbit interaction is given by $V_{SO} = \frac{1}{\hbar^2} V_{so}(r) \hat{\vec{l}} \cdot \hat{\vec{s}}$

- We can calculate the dot product

$$\langle \hat{\vec{l}} \cdot \hat{\vec{s}} \rangle = \frac{1}{2} (\hat{j}^2 - \hat{l}^2 - \hat{s}^2) = \frac{\hbar^2}{2} [j(j+1) - l(l+1) - \frac{3}{4}]$$

- Because of the addition rules, $j = l \pm \frac{1}{2}$

$$\langle \hat{\vec{l}} \cdot \hat{\vec{s}} \rangle = \begin{cases} l \frac{\hbar^2}{2} & \text{for } j=l+\frac{1}{2} \\ -(l+1) \frac{\hbar^2}{2} & \text{for } j=l-\frac{1}{2} \end{cases}$$

Spin-Orbit Coupling

- when the spin is **aligned** with the angular momentum $j = l + \frac{1}{2}$ the potential becomes more negative, i.e. the well is deeper and the state more tightly bound.

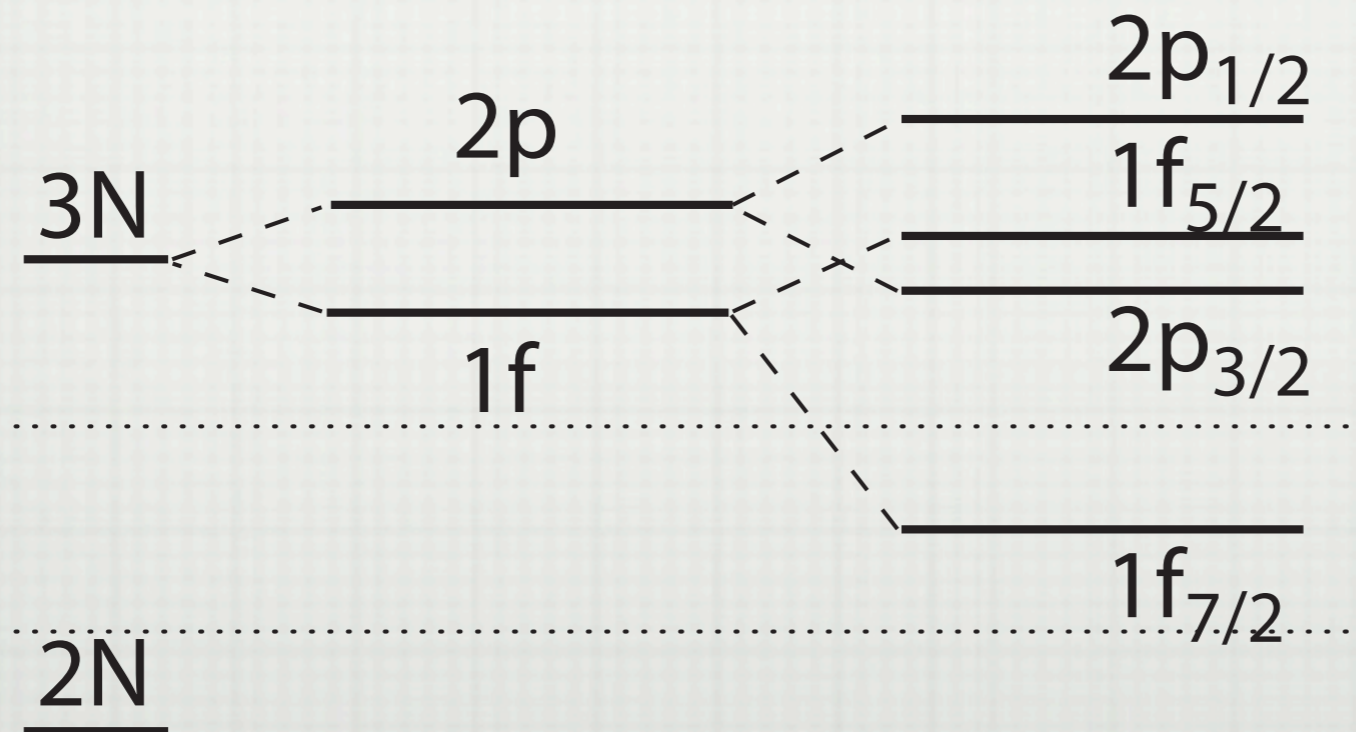
- when spin and angular momentum are **anti-aligned** $j = l - \frac{1}{2}$ the system's energy is higher.

- The difference in energy is
$$\Delta E = \frac{V_{so}}{2} (2l + 1)$$

Thus it increases with l .

Example

- 3N level, with $l=3$ (1f level) $j=7/2$ or $j=5/2$
 - Level is pushed so down that it forms its own shell

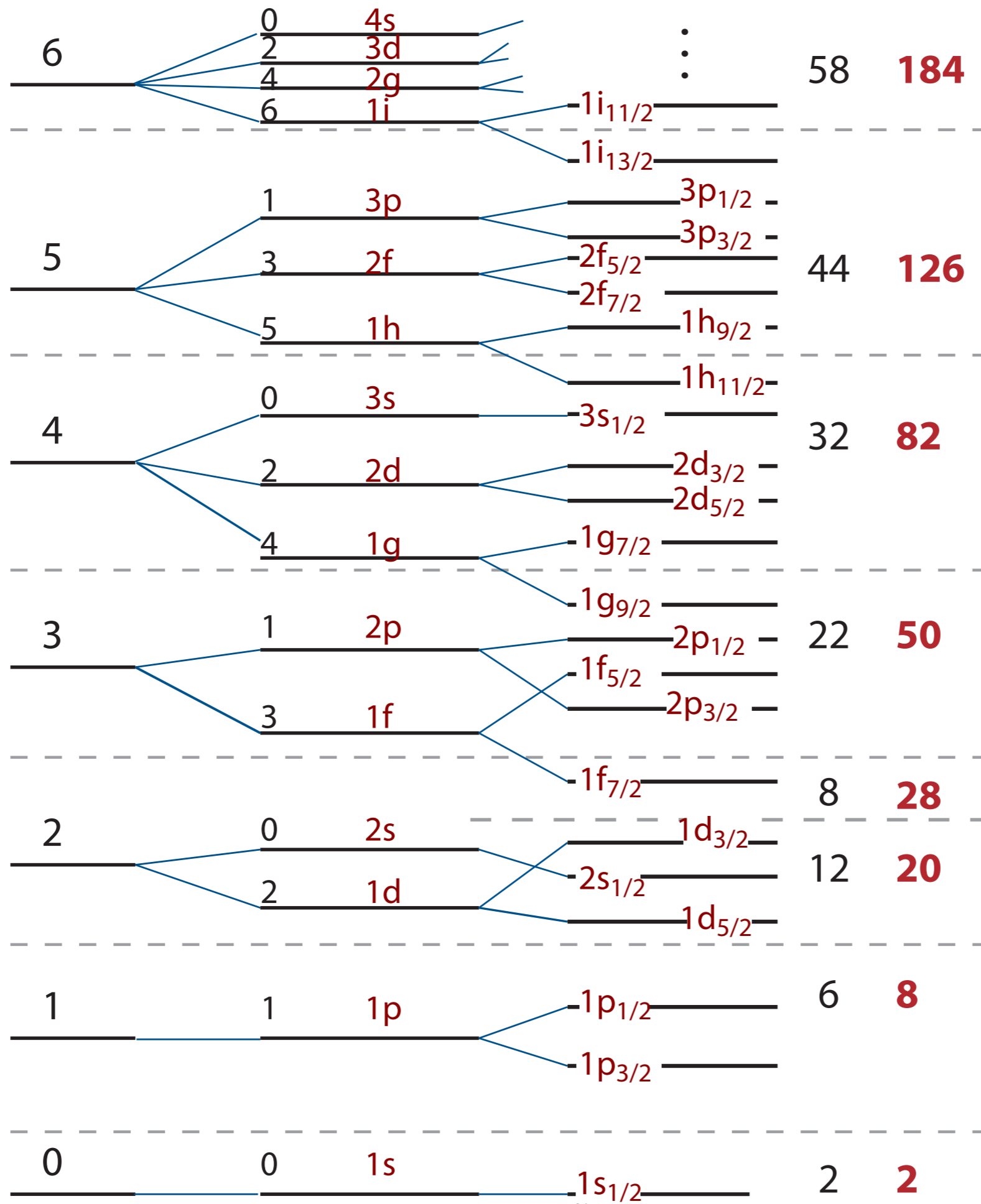


Harmonic Oscillator

Harmonic oscillator with angular momentum

Spin-Orbit Potential

N	ℓ	Spec. Not	\mathcal{D}	M.N.	Spec. Not	Spin-orbit	\mathcal{D}	M.N.
6	0	4s	56	168	4s	1j _{15/2}	16	184
	2	3d			3d			
	4	2g			2g			
	6	1i			1i			
5	1	3p	42	112	3p	1i _{13/2}	14	126
	3	2f			2f			
	5	1h			1h			
4	0	3s	30	70	3s	1h _{11/2}	12	82
	2	2d			2d			
	4	1g			1g			
3	1	2p	20	40	2p	1g _{9/2}	10	50
	3	1f			1f			
2	0	2s	12	20	2s	1f _{7/2}	8	28
	2	1d			1d			
1	1	1p	6	8	1p		6	8
0	0	1s	2	2	1s		2	2



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