

# SINUSOIDAL RADIATION BY ANTENNAS

Poisson equations:

$$\nabla^2 \Phi = -\frac{\rho}{\epsilon}$$

$$\nabla^2 \bar{A} = -\mu_0 \bar{J} \quad (\bar{B} \triangleq \nabla \times \bar{A})$$

Retarded potentials:

$$\Phi_p = \frac{1}{4\pi\epsilon_0} \int_{V_q} \frac{\rho_q(t - r_{pq}/c)}{r_{pq}} dV_q$$

$$\bar{A}_p = \frac{\mu_0}{4\pi} \int_{V_q} \frac{\bar{J}_q(t - r_{pq}/c)}{r_{pq}} dV_q$$

$$\underline{\Phi}(\omega) = \frac{1}{4\pi\epsilon_0} \int_{V_q} \frac{\underline{\rho}(\omega)}{r_{pq}} e^{-jk r_{pq}} dV_q$$

$$\bar{A}(\omega) = \frac{\mu_0}{4\pi} \int_{V_q} \frac{\bar{J}(\omega)}{r_{pq}} e^{-jk r_{pq}} dV_q$$

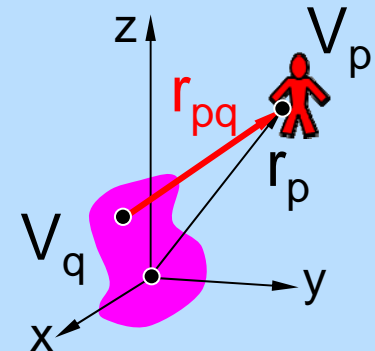
Radiation  $\cong$  static solution if  $kr_{pq} = \frac{2\pi r_{pq}}{\lambda} \ll 1$ , or  $r_{pq} \ll \frac{\lambda}{2\pi}$

Finding fields from potentials:

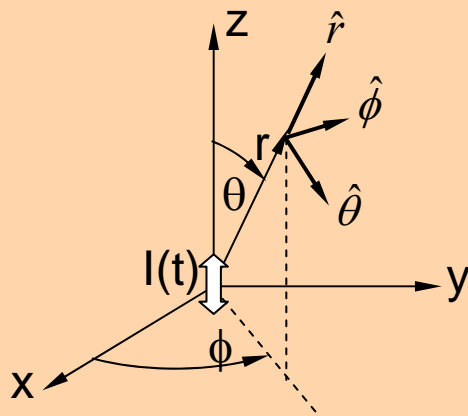
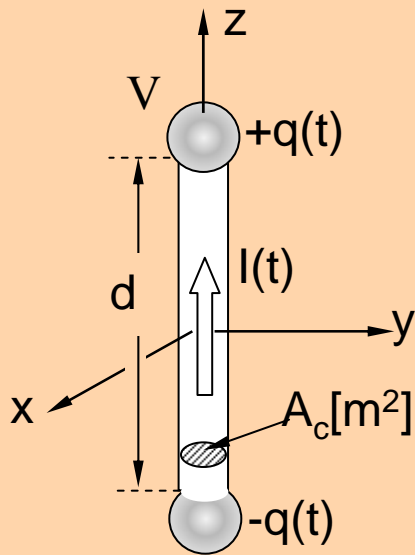
$$\bar{B} = \nabla \times \bar{A}, \quad \bar{E} = -\nabla \Phi - j\omega \bar{A} = -\nabla \times \frac{\bar{H}}{j\omega\epsilon}$$

$$k = \omega \sqrt{\mu_0 \epsilon_0} = \frac{2\pi}{\lambda_0}$$

Algorithm for general problem:  $\{\underline{\rho}, \bar{J}\} \Rightarrow \{\Phi, \bar{A}\} \Rightarrow \{\bar{E}, \bar{H}\}$



# HERTZIAN DIPOLE RADIATOR



$$\begin{aligned} \bar{\underline{A}}_p(\omega) &= \frac{\mu_0}{4\pi} \int_{V_q} \frac{\bar{\underline{J}}_q(\omega)}{r_{pq}} e^{-jkr_{pq}} dV_q \\ &\cong \underbrace{(\hat{r} \cos \theta - \hat{\theta} \sin \theta)}_{\hat{z}} \frac{\mu_0 I(\omega) d}{4\pi r_{pq}} e^{-jkr_{pq}} \end{aligned}$$

Where:  $\bar{\underline{J}}_q = \hat{z} \underline{J}_q$  for  $-\frac{d}{2} < z < \frac{d}{2}$

And provided that  $d \ll r$   
(so  $r$  can come outside the integral)

$$\bar{\underline{H}} = \frac{\nabla \times \bar{\underline{A}}}{\mu_0} = \frac{1}{\mu_0 r^2 \sin \theta} \det \begin{vmatrix} \hat{r} & r\hat{\theta} & r \sin \theta \hat{\phi} \\ \partial/\partial r & \partial/\partial \theta & \partial/\partial \phi \\ A_r & rA_\theta & r \sin \theta A_\phi \end{vmatrix}$$

# SOLVING FOR HERZIAN DIPOLE FIELDS

$$\bar{\mathbf{H}} = \frac{\nabla \times \bar{\mathbf{A}}}{\mu_0} = \hat{\phi} \frac{jk \underline{I} d}{4\pi r} e^{-jkr} \left[ 1 + \frac{1}{jkr} \right] \sin \theta$$

$$\begin{aligned} \bar{\mathbf{E}} &= \frac{\nabla \times \bar{\mathbf{H}}}{j\omega\epsilon_0} \\ &= j \frac{k \underline{I} d \eta_0}{4\pi r} e^{-jkr} \left\{ \hat{r} \left[ \frac{1}{jkr} + \frac{1}{(jkr)^2} \right] 2 \cos \theta + \hat{\theta} \left[ 1 + \frac{1}{jkr} + \frac{1}{(jkr)^2} \right] \sin \theta \right\} \end{aligned}$$

Far field:  $r \gg \frac{\lambda}{2\pi}$

$$\bar{\mathbf{H}}_{\text{ff}} \cong \hat{\phi} \frac{jk \underline{I} d}{4\pi r} e^{-jkr} \sin \theta$$

$$\bar{\mathbf{E}}_{\text{ff}} \cong \hat{\theta} \frac{jk \underline{I} d \eta_0}{4\pi r} e^{-jkr} \sin \theta$$

Near field:  $r \ll \frac{\lambda}{2\pi}$

$$\bar{\mathbf{H}}_{\text{nf}} \cong \hat{\phi} \frac{\underline{I} d}{4\pi r^2} e^{-jkr} \sin \theta$$

(yields Biot-Savart law)

$$\bar{\mathbf{E}}_{\text{nf}} \cong -j \frac{\underline{I} d \eta_0}{4\pi k r^3} e^{-jkr} (\hat{r} 2 \cos \theta + \hat{\theta} \sin \theta)$$

$$\bar{\mathbf{H}}(\bar{\mathbf{r}}, t) = \iiint_{V'} \frac{\hat{\phi} \underline{J} \sin \theta}{4\pi |\bar{\mathbf{r}} - \bar{\mathbf{r}}'|^2} dv. \quad \text{Since } \hat{\phi} \underline{J} \sin \theta = \underline{\mathbf{J}} \times \frac{(\bar{\mathbf{r}} - \bar{\mathbf{r}}')}{|\bar{\mathbf{r}} - \bar{\mathbf{r}}'|}, \quad \bar{\mathbf{H}}(\bar{\mathbf{r}}, t) = \iiint_{V'} \frac{\underline{\mathbf{J}} \times (\bar{\mathbf{r}} - \bar{\mathbf{r}}')}{4\pi |\bar{\mathbf{r}} - \bar{\mathbf{r}}'|^3} dv$$

# NEAR AND FAR FIELDS OF HERTZIAN DIPOLE

Far field:  $r \gg \frac{\lambda}{2\pi}$  (radiation)

$$\underline{\bar{E}}_{ff} \cong \hat{\theta} \frac{jkI_d \eta_0}{4\pi r} e^{-jkr} \sin \theta$$

$$\underline{\bar{H}}_{ff} \cong \hat{\phi} \frac{jkI_d}{4\pi r} e^{-jkr} \sin \theta$$

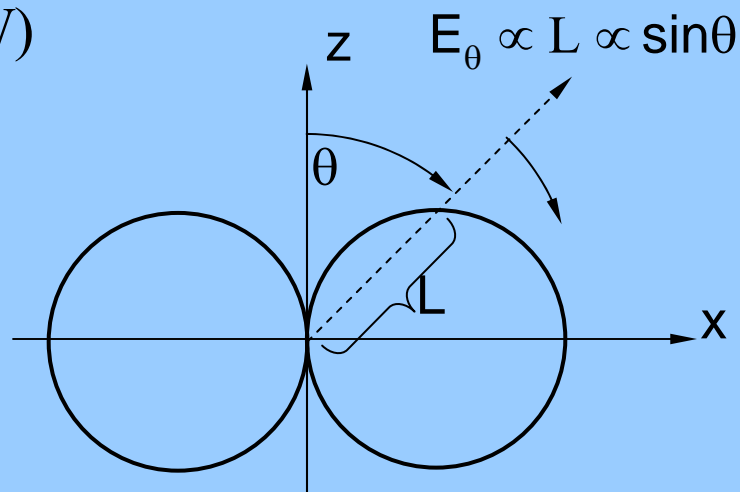
$$\underline{\bar{E}}_{ff} = \eta_0 \underline{\bar{H}}_{ff} \times \hat{r} \quad (\text{UPW})$$

$$(\eta_0 = \sqrt{\frac{\mu_0}{\epsilon_0}} = 377\Omega)$$

Near field:  $r \ll \frac{\lambda}{2\pi}$  (quasistatic)

$$\underline{\bar{E}}_{nf} \cong -\frac{jI_d \eta_0}{4\pi k r^3} e^{-jkr} (\hat{r} 2 \cos \theta + \hat{\theta} \sin \theta)$$

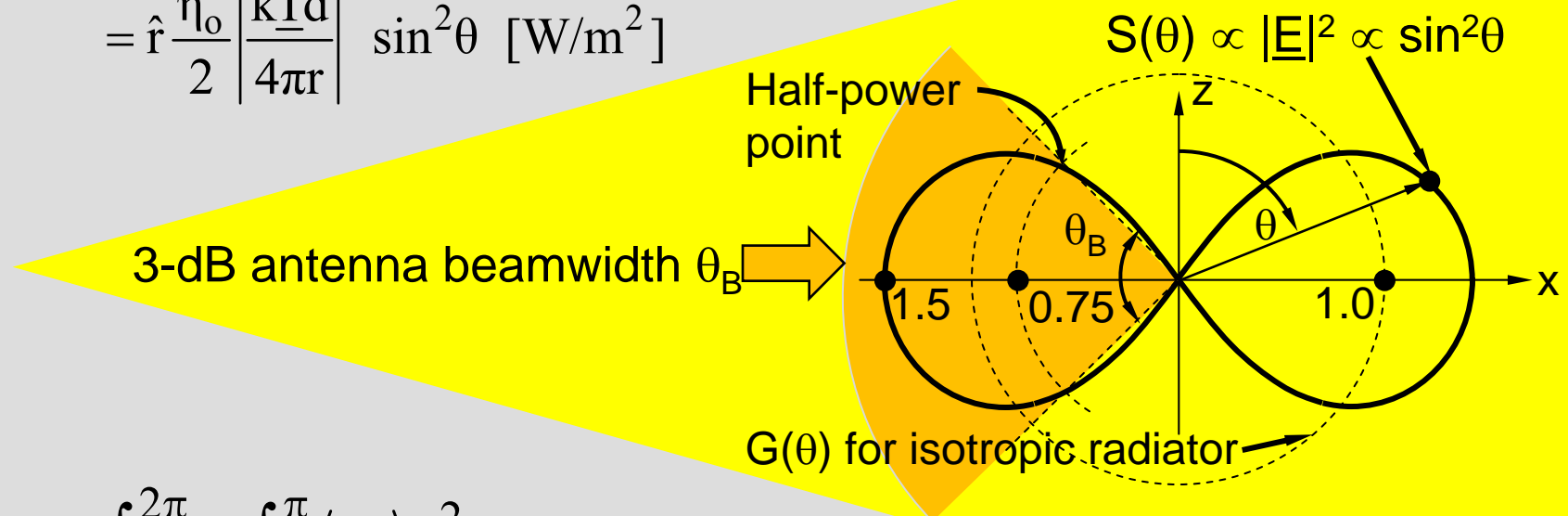
$$\underline{\bar{H}}_{nf} \cong \hat{\phi} \frac{I_d}{4\pi r^2} e^{-jkr} \sin \theta$$



# POWER RADIATED BY HERTZIAN DIPOLE

$$\langle \bar{S}(t) \rangle = \frac{1}{2} \text{Re} \{ \bar{S} \} = \frac{1}{2} \text{Re} \{ \bar{E} \times \bar{H}^* \} = \hat{r} \frac{1}{2\eta_0} |\bar{E}_0|^2$$

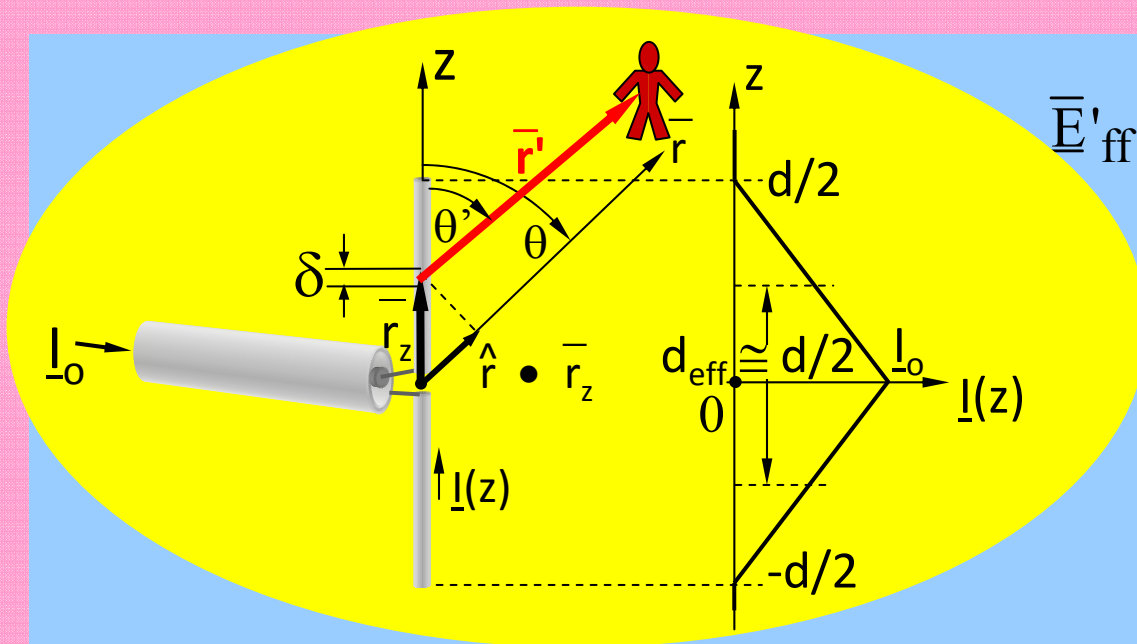
$$= \hat{r} \frac{\eta_0}{2} \left| \frac{k \underline{I} d}{4\pi r} \right|^2 \sin^2 \theta \quad [\text{W/m}^2]$$



$$P_R = \int_0^{2\pi} d\phi \int_0^\pi \langle S_r \rangle r^2 \sin\theta \, d\theta$$

$$= \pi \eta_0 \left| \frac{k \underline{I} d}{4\pi} \right|^2 \int_0^\pi \sin^3 \theta \, d\theta = \frac{\eta_0}{12\pi} |k \underline{I} d|^2 \cong 395 \left| \frac{\underline{I} d}{\lambda} \right|^2 \quad [\text{W}]$$

# SHORT DIPOLE ANTENNAS



$$\underline{\bar{E}}'_{ff} \cong \hat{\theta}' \frac{jkI(z)\delta\eta_0}{4\pi r'} e^{-jkr'} \sin\theta'$$

Wire  $\cong$  TEM line  
open-circuited,  
so current is a  
truncated sinusoid  
 $\Rightarrow$  triangle.

$$\begin{aligned} \underline{\bar{E}}_{ff} &= \int_{-d/2}^{d/2} \underline{\bar{E}}'_{ff}(r', \theta') dz = \frac{jk\eta_0}{4\pi} \int_{-d/2}^{d/2} \hat{\theta}' \frac{\sin\theta'}{r'} I(z) e^{-jkr'} dz \\ &\cong \hat{\theta} \sin\theta \frac{jk\eta_0}{4\pi r} e^{-jkr} \int_{-d/2}^{d/2} I(z) dz \cong \hat{\theta} \sin\theta \frac{jk\eta_0}{4\pi r} I_0 d_{eff} e^{-jkr} \end{aligned}$$

Has same far fields as does a Hertzian dipole, except  $I_d \rightarrow I_0 d_{eff}$

# ANTENNA RADIATION RESISTANCE

$$P_T = \frac{\eta_0}{12\pi} |k \underline{I}_0 d_{\text{eff}}|^2 \cong 395 \left| \frac{\underline{I}_0 d_{\text{eff}}}{\lambda} \right|^2 \text{ [W]}$$

$$= \frac{1}{2} |\underline{I}_0|^2 R_r \Rightarrow \text{Radiation resistance } R_r = \frac{2P_T}{|\underline{I}_0|^2} = \frac{2\eta_0\pi}{3} \left( \frac{d_{\text{eff}}}{\lambda} \right)^2 \text{ ohms}$$

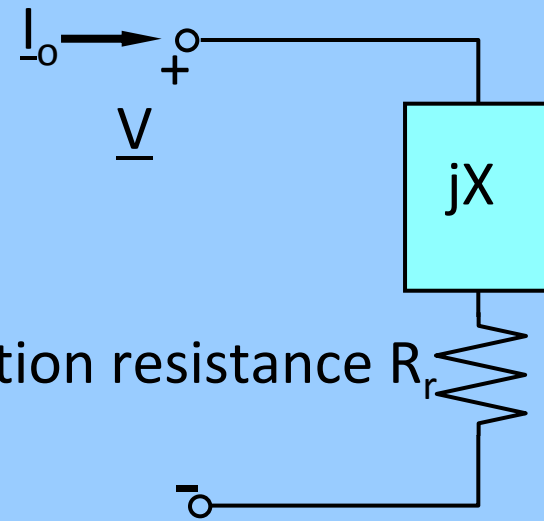
For a triangular current distribution,  
center fed,  $d_{\text{eff}} \cong d/2$ .

Short dipole antenna example:

AM radio at 1 MHz has  $\lambda = 300$  m.

Car antenna 2m long has:

$$R_r = \frac{2 \cdot 377 \cdot 3.14}{3} \left( \frac{1}{300} \right)^2 \cong 0.02 \text{ ohms}$$



Radiation resistance  $R_r$

$jX$  can similarly be found by measuring or computing  $\underline{S}$  into an antenna; it results from net electric or magnetic energy in the near field.

# ANTENNA GAIN OVER ISOTROPIC

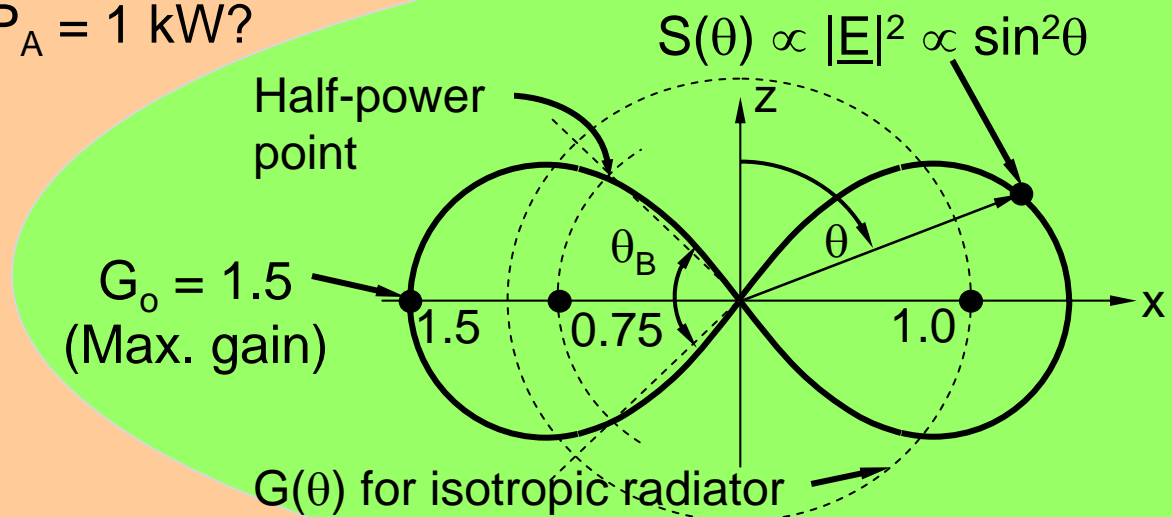
$$G(\theta, \phi) \equiv \frac{P(r, \theta, \phi)}{(P_A/4\pi r^2)} = \frac{|\underline{E}(r, \theta, \phi)|^2 / 2\eta_0}{(P_A/4\pi r^2)} = \frac{|\sin \theta \frac{jk\eta_0}{4\pi r} I_0 d_{\text{eff}}|^2 / 2\eta_0}{\frac{\eta_0}{12\pi} |k I_0 d_{\text{eff}}|^2 / 4\pi r^2} = \frac{3}{2} \sin^2 \theta$$

Example: What maximum  $\bar{E}$  [V/m] is produced at 10 km by a short dipole radiating  $P_A = 1$  kW?

Solution:

$$\begin{aligned} \frac{|\bar{E}(r, \theta, \phi)|^2}{2\eta_0} &= P(r, \theta, \phi) \\ &= \frac{(1.5 \sin^2 \theta) P_A}{4\pi r^2} \end{aligned}$$

$$\Rightarrow |\bar{E}(r, \theta, \phi)| = \sqrt{(1.5 \sin^2 \theta) 10^3 \eta_0 / 2\pi 10^8} = 30 \text{ mV/m}$$





# ANTENNA DIRECTIVITY AND RADIATION EFFICIENCY

**Gain over isotropic**

$$G(\theta, \phi) \equiv \frac{P(r, \theta, \phi)}{(P_A/4\pi r^2)}$$

**Antenna directivity**

$$D(\theta, \phi) \equiv \frac{P(r, \theta, \phi)}{(P_T/4\pi r^2)}$$

$$\frac{G(\theta, \phi)}{D(\theta, \phi)} = \frac{P_T}{P_A} = \frac{\text{Power transmitted}}{\text{Power available}} \triangleq \eta_r \quad \text{"radiation efficiency"} \leq 1$$

$$\oint_{4\pi} D(\theta, \phi) \sin \theta \, d\theta \, d\phi = 4\pi \quad \text{since} \quad \oint_{4\pi} P(r, \theta, \phi) r^2 \sin \theta \, d\theta \, d\phi = P_T$$

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