

Fourier analysis of discretized PDEs

PDE

$$\frac{\partial U}{\partial t} + c \frac{\partial U}{\partial x} = 0$$

Periodic: $U(x+nL, t) = U(x, t)$
for any integer n

Substitute a Fourier mode:

$$U(x, t) = \hat{U}_m(t) \exp(ik_m x)$$

$$k_m = \frac{2\pi m}{L} \text{ for any integer } m$$

$$\begin{aligned} \frac{d\hat{U}_m}{dt} \exp(ik_m x) + \\ ik_m c \hat{U}_m(t) \exp(ik_m x) \\ = 0 \end{aligned}$$

Semi-discrete w/central space

$$\frac{dV_j}{dt} + c \frac{V_{j+1} - V_{j-1}}{2\Delta x} = 0$$

Periodic: $V_{j+n\frac{L}{\Delta x}}(t) = V_j(t)$
for any integer n

Substitute a Fourier mode:

$$\begin{aligned} V_j(t) &= \hat{V}_m(t) \exp(ik_m x_j) \\ &= \hat{V}_m(t) \exp(ik_m j \Delta x) \end{aligned}$$

$$\begin{aligned} \frac{d\hat{V}_m}{dt} \exp(ik_m j \Delta x) + \\ \frac{c}{2\Delta x} \hat{V}_m(t) \left[\exp(ik_m (j+1) \Delta x) - \exp(ik_m (j-1) \Delta x) \right] \\ = 0 \end{aligned}$$

$$\frac{d\hat{U}_m}{dt} = \underbrace{-ik_m c}_{\lambda_u} \hat{U}_m$$

$$\frac{d\hat{V}_m}{dt} = -\frac{c}{2\Delta x} [\exp(ik_m \Delta x) - \exp(-ik_m \Delta x)] \hat{V}_m$$

$$\frac{d\hat{V}_m}{dt} = \underbrace{-i \frac{c}{\Delta x} \sin(k_m \Delta x)}_{\lambda_v} \hat{V}_m$$

Comment: For $k_m \Delta x \ll 1$ (wave has long wavelength compared to Δx)

$$\lambda_v \approx -i \frac{c}{\Delta x} (k_m \Delta x) = -i c k_m = \lambda_u$$

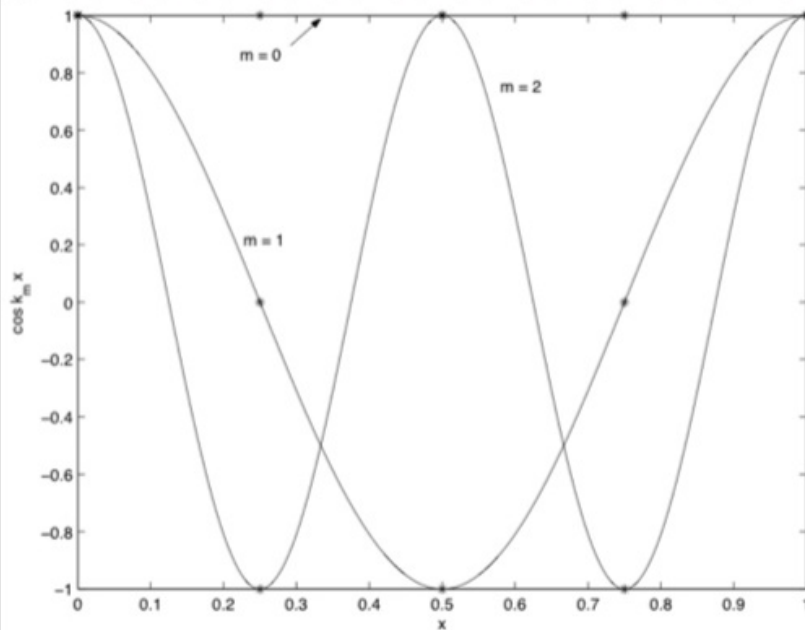
Discrete eigenvalues converge to PDE eigenvalues as $\Delta x \rightarrow 0$.

How does $e^{ik_m x}$ behave? $k_m = \frac{2\pi m}{L} \Rightarrow \lambda_m = \frac{2\pi}{|k_m|} = \frac{L}{|m|}$

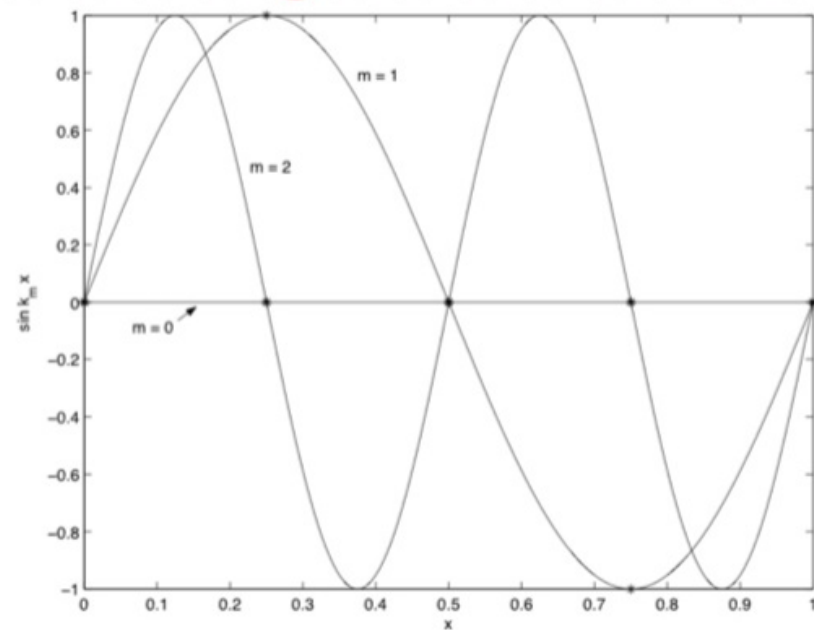
Example: $L=1, \Delta x = 1/4$

Note: $\frac{\lambda_m}{\Delta x} = \frac{L}{|m|\Delta x}$

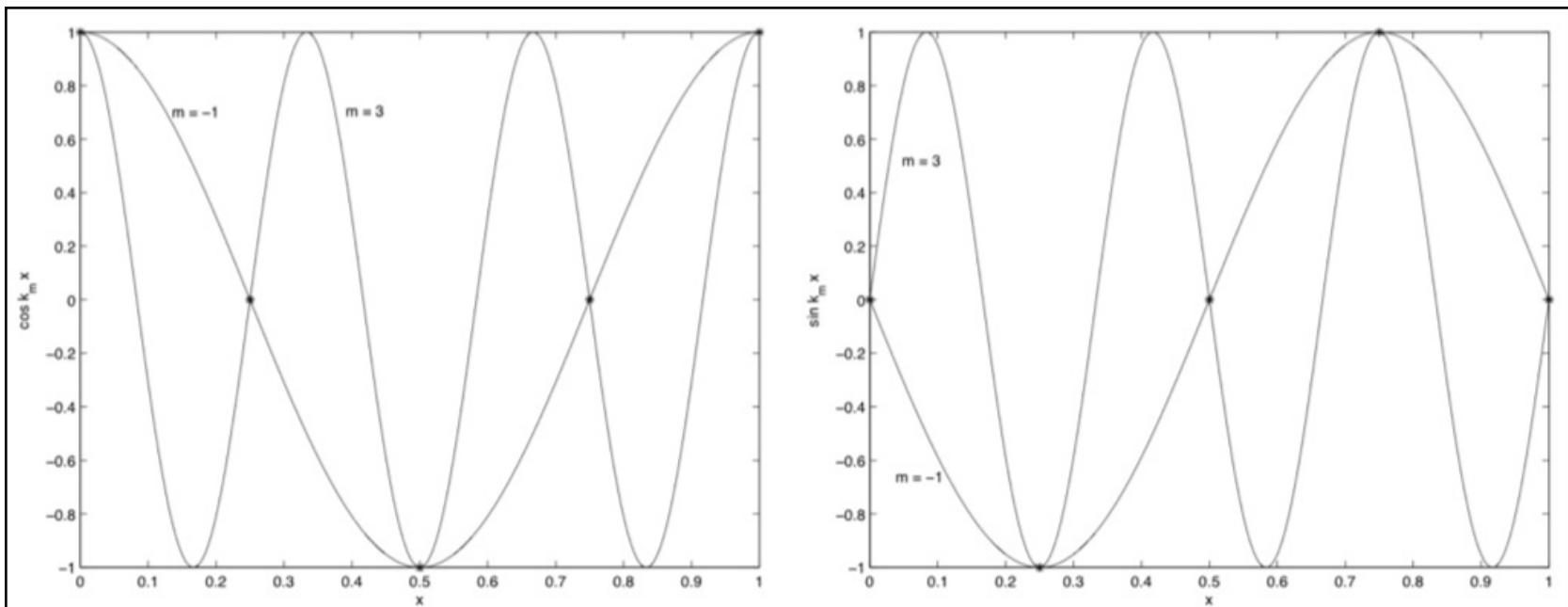
When $\lambda_m < 2\Delta x$, aliasing!
 Equivalently when $|k_m|\Delta x > \pi$
 aliasing occurs!



(a) Real $e^{ik_m x} = \cos k_m x$



(b) Imag $e^{ik_m x} = \sin k_m x$

(a) Real $e^{ik_m x} = \cos k_m x$ (b) Imag $e^{ik_m x} = \sin k_m x$

$$-\pi \leq \beta_m \equiv k_m \Delta x \leq \pi$$

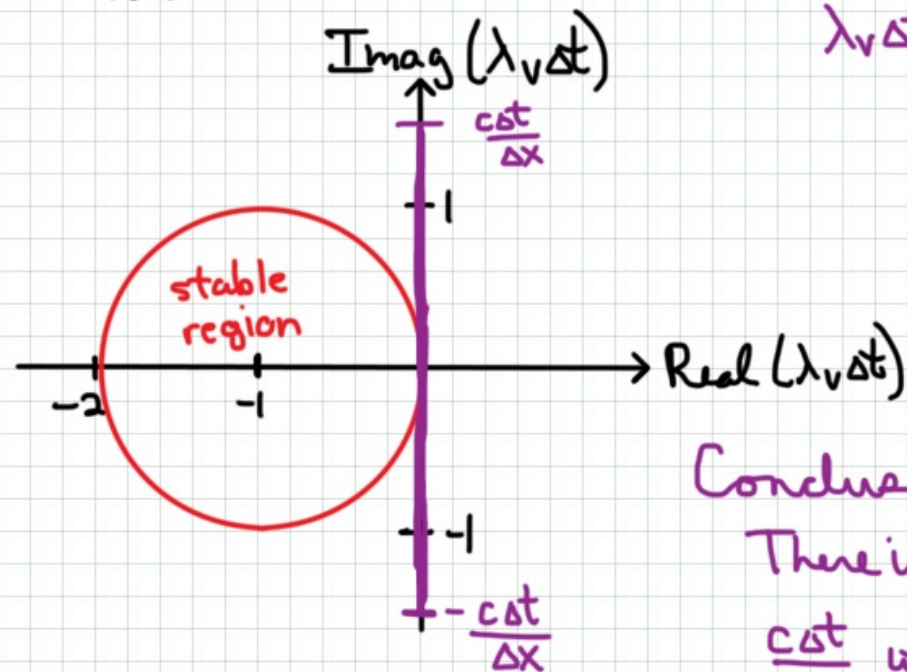
Now, let's return to our semi-discrete analysis:

$$\frac{d\hat{V}_m}{dt} = \lambda_v \hat{V}_m \quad \text{where } \lambda_v = -i \frac{c}{\Delta x} \sin \beta_m$$

Determine stability for a specific time integration method:

Criterion: λ_{vst} must be in stability region for chosen time integration.

Example: Forward Euler:



$$\lambda_{vst} = -i \frac{\cot}{\Delta x} \sin \beta_m$$

Conclusion:

There is no value of $\frac{\cot}{\Delta x}$ which will be stable.

Summary of Fourier analysis of stability:

(1) Substitute $V_j(t) = \hat{V}_m(t) \exp(ij\beta_m)$ into semi-discrete method

(2) Determine eigenvalue $\lambda_V(\beta_m)$ such that $\frac{d\hat{V}_m}{dt} = \lambda_V \hat{V}_m$

(3) Determine timestep requirements for all $\lambda_V(\beta_m)\Delta t$ to be inside stable region for a chosen time integrator

Example: Diffusion: $\frac{\partial U}{\partial t} = \mu \frac{\partial^2 U}{\partial x^2}$ for $\mu > 0$

Spatial discretization: centered, second order

Time integrator: Forward Euler

Determine Δt requirements for stability

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